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Numerical Simulation of Shear Flows Over Wavy Boundaries

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Numerical Simulation of Shear Flows Over Wavy Boundaries

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SYMBOLS

a	Amplitude of wall displacement
A	Fourier-Chebyshev amplitude of v ; see (16)
C_p	Pressure coefficient
f	Wall shape function
G_k	See (20)
H	$H(Y) = dY/d\eta$; see (17)
k	Wavenumber of wall displacement
\bar{k}	Turbulent kinetic energy
m	X-wavenumber
N	Number of retained modes in x_1 or ζ
p	Pressure
Re_x	Ux/ν , Reynolds number
t	Time
T	Reynolds stress
T_n	Chebyshev polynomial of degree n
u_k	Eigenfunction; see (18)
u	Velocity in x-direction
U	Free-stream velocity in x-direction
v	Velocity in y-direction
V_k	See (19)
x	Unmapped Cartesian coordinate
x_1	$x_1 = x$
x_2	$x_2 = y$
X	Mapped coordinate; see (14) or (15)
y	Unmapped Cartesian coordinate
Y	Mapped coordinate; see (14) or (15)
α	$= 1,2$; coordinate label
Δt	Time step
Δx	Space discretization interval
Δu	Velocity fluctuation
δ	Boundary layer thickness
η	Coordinate in conformal mapped system
ϕ_0	Phase of wall displacement

λ_k	Eigenvalue; see (18)
μ_T	Eddy viscosity coefficient
ν	Viscosity
π	Pi
ρ	Density
$\bar{\rho}$	Parameter used in (17)
σ	Parameter used in viscous stabilization; see (13)
ζ	Coordinate in conformally mapped system

Superscripts

\sim	Convective update
\wedge	Pressure update
$—$	Viscous update
\equiv	Viscous stabilization
$'$	Derivative with respect to $\zeta + i\eta$; Also, fluctuation from mean flow

I. INTRODUCTION

The study of mechanisms of possible drag reduction effects in flows over wavy walls as well as the study of mechanisms of generation of surface waves by wind requires understanding of the detailed flows over wavy surfaces. In this paper, we describe a computer code based on spectral methods to study two-dimensional incompressible flows in wavy geometries.

Work on flows in wavy geometries began with the classical analysis of Kelvin and Helmholtz of linearized inviscid surface waves¹. In the absence of mean-flow-shear effects, inviscid theory predicts that, for wall displacements of the form $y = a \cos(kx + \phi_0)$ with $ka \ll 1$ (so that the problem may be linearized in terms of ka), the pressure distribution is 180° out of phase with the wall displacement and the pressure coefficient $C_p = |p|_{\max} / \frac{1}{2} \rho U^2$ has the value $2ka$ at the wall. Miles² extended the inviscid theory of account for shear in the mean flow, still restricting attention to very gentle waves. Miles' theory improves the prediction of the pressure coefficients. However, because of the inviscid nature of the analysis, the pressure phase shift at the wall is still 180° .

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Benjamin³ analyzed high Reynolds number laminar flows over wavy walls including the effects of shear still assuming $ka \ll 1$. Benjamin's theory typically predicts a pressure phase shift of 210° relative to the displacement of a solid wavy boundary. The theory of Miles and Benjamin has also been extended to turbulent flows. Davis⁴ analyzes turbulent flow over wavy surfaces by postulating that Reynolds stresses are, to a first approximation, constant along lines of constant η , where (ζ, η) is a first-order accurate orthogonal coordinate system with the wavy boundary at $\eta = 0$. Davis' analysis is, therefore, also limited to small ka . Markatos⁵ has recently presented numerical results for transfer of heat, mass and momentum in flow over an evaporating wavy water surface. While Markatos' numerical results are at best first-order accurate in space, they go a long way towards providing understanding of the nature of flow over wavy walls.

The work described in this paper extends previous work in several ways. The full time-dependent two-dimensional Navier-Stokes equations are solved using spectral methods to achieve high spatial accuracy and high-order time-splitting methods together with conformal mapping methods to allow simulation of flow over steep waves. Results for laminar flows at

Reynolds numbers up to $Re_x = 4 \times 10^5$ are presented in Sec. III. Turbulent flows over wavy walls are simulated by using a spatially varying eddy viscosity distribution. Some preliminary results and comparisons with experiment are given in Sec. IV.

II. METHODS OF SOLUTION

The two-dimensional Navier-Stokes equations are

$$\frac{\partial v_\alpha}{\partial t} + \sum_{\beta=1}^2 v_\beta \frac{\partial v_\alpha}{\partial x_\beta} = - \frac{\partial p}{\partial x_\alpha} + \sum_{\beta=1}^2 \frac{\partial}{\partial x_\beta} T_{\alpha\beta} \quad (1)$$

with $\alpha = 1, 2$,

$$\sum_{\alpha=1}^2 \frac{\partial v_\alpha}{\partial x_\alpha} = 0, \quad (2)$$

where the stress tensor may include both laminar (viscous) stresses, $\nu \nabla^2 v_\alpha$, and turbulent Reynolds stresses (to be discussed briefly in Sec. IV below). Eqs. (1) - (2) are to be solved in the region

$$0 \leq x_1 < 2\pi, \quad f(x_1, t) \leq x_2 < \infty \quad (3)$$

above the wall $x_2 = f(x_1, t)$. In the present paper, we restrict attention to surfaces of the form

$$f(x_1, t) = a \cos kx \quad (4)$$

independent of t with periodic inflow-outflow boundary conditions in x_1 , viz.

$$v_{\alpha}(x_1 + 2\pi, x_2, t) = v_{\alpha}(x_1, x_2, t) \quad (5)$$

Extensions to study the effects of time-dependent geometries and inflow-outflow boundary conditions are presently being made.

A conformal mapping technique⁶ is used to transform the region (4) into the region

$$0 \leq \zeta \leq 2\pi, \quad 0 \leq \eta < \infty$$

If $a \leq 1$, the conformal mapping coefficients are accurately generated in only order $N \log N$ operations where N is the resolution along $x_1(\zeta)$, so time dependent geometries can be efficiently handled. A further (non-conformal) stretching transformation of η is used to implement the spectral methods described below see Eqs. (14)-(15)

Eqs. (1) - (2) are solved numerically by a fractional step procedure. The fractional steps are: (i) Convective update

An intermediate field \tilde{v}_{α} is found so that

$$\tilde{v}_{\alpha} = v_{\alpha}(t) + \frac{3}{2} \Delta t N_{\alpha}(t) - \frac{1}{2} N_{\alpha}(t - \Delta t) \Delta t \quad (6)$$

where

$$N_{\alpha} = \sum_{\beta} v_{\beta} \frac{\partial}{\partial v_{\beta}} v_{\alpha}$$

Here a second-order accurate Adams-Bashforth scheme is applied.

(ii) Pressure update

$$\hat{v}_\alpha = \tilde{v}_\alpha - \frac{\partial \zeta}{\partial x_\alpha} \frac{\partial p}{\partial \zeta} - \frac{\partial \eta}{\partial x_\alpha} \frac{\partial p}{\partial \eta} \quad (7)$$

$$\sum_{\alpha=1}^2 \left(\frac{\partial \zeta}{\partial x_\alpha} \frac{\partial \hat{v}_\alpha}{\partial \zeta} + \frac{\partial \eta}{\partial x_\alpha} \frac{\partial \hat{v}_\alpha}{\partial \eta} \right) = 0 \quad (8)$$

with the boundary condition

$$\sum_{\alpha=1}^2 \frac{\partial \eta}{\partial x_\alpha} \hat{v}_\alpha = 0 \quad \text{at } \eta = 0. \quad (9)$$

Substituting (7) into (8) gives the Poisson equation

$$\frac{\partial^2 p}{\partial \zeta^2} + \frac{\partial^2 p}{\partial \eta^2} = \frac{1}{|F'|^2} \sum_{\alpha=1}^2 \left(\frac{\partial \zeta}{\partial x_\alpha} \frac{\partial \tilde{v}_\alpha}{\partial \zeta} + \frac{\partial \eta}{\partial x_\alpha} \frac{\partial \tilde{v}_\alpha}{\partial \eta} \right) \quad (10)$$

with

$$|F'|^2 = \sum_{\alpha=1}^2 \left(\frac{\partial \zeta}{\partial x_\alpha} \right)^2 = \sum_{\alpha=1}^2 \left(\frac{\partial \eta}{\partial x_\alpha} \right)^2, \quad (11)$$

where the conformal nature of the map from (x_1, x_2) to (ζ, η) is used.

(iii) Viscous update

$$\bar{v}_\alpha = \hat{v}_\alpha + \Delta t \sum_{\alpha=1}^2 \frac{\partial}{\partial x_\beta} T_{\alpha\beta}(t) \quad (12)$$

with the boundary conditions

$$\bar{v}_\alpha = 0 \quad \text{at } \eta = 0.$$

This Euler step is only first-order accurate but will be improved upon later. The time step restrictions for numerical stability of (12) are particularly severe, especially in the case of turbulent flows, so the next fractional step is designed to stabilize the scheme.

(iv) Viscous stabilization

$$\bar{\bar{v}}_\alpha = \frac{\sigma}{|F'|_{\min}^2} \nabla_\zeta^2 \bar{v}_\alpha = \bar{v}_\alpha - \frac{\sigma}{|F'|_{\min}^2} \nabla_\zeta^2 v_\alpha(t) \quad (13)$$

with

$$\bar{\bar{v}}_\alpha = 0 \quad \text{at } \eta = 0,$$

where $\nabla_\zeta^2 = \partial^2/\partial \zeta^2 + \partial^2/\partial \eta^2$ is the Laplacian with respect to ζ, η . Here σ is a parameter chosen to achieve optimal stabilization and accuracy.

(v) Local extrapolation

For time dependent problems requiring accurate calculation of transient effects, a local extrapolation procedure⁷ is used to achieve global second-order accuracy in time. Further details and caveats on the splitting procedure described above are given in Ref. 7.

In order to implement this algorithm in a spectral method⁸, the semi-infinite region $0 \leq \zeta \leq 2\pi$, $0 \leq \eta < \infty$ is mapped into the finite rectangular domain

$$0 \leq X \leq 2\pi, \quad -1 \leq Y < 1$$

by

$$X = \zeta, \quad Y = \frac{2\eta}{\eta+L} - 1 \quad (14)$$

or

$$X = \zeta, \quad Y = 1 - 2e^{-\eta/L} \quad (15)$$

for some suitable mapping scale L .

Next, all dependent variables are expanded in Fourier-Chebyshev series of the form

$$v_\alpha(X, Y, t) = \sum_{n=0}^N \sum_{|m| < M} A_\alpha(n, m, t) e^{imX} T_n(Y) \quad (16)$$

where $T_n(Y) = \cos(n \cos^{-1} Y)$ is the Chebyshev polynomial of degree n .

Technical details regarding spectral methods and their implementation are given in Ref. 8 and references given therein. For illustration purposes, the method of solution of the Poisson equation (10) or (13) will now be described. After the transformation (14) or (15), the equation to be solved is of the form

$$\bar{\rho} u + \frac{\partial^2 u}{\partial X^2} + H(Y) \frac{\partial H(Y)}{\partial Y} \frac{\partial u}{\partial Y} = g(X, Y). \quad (17)$$

Introducing the eigenvalue problem

$$H(Y) \frac{\partial}{\partial Y} H(Y) \frac{\partial u_k}{\partial Y} = \lambda_k u_k \quad (18)$$

subject to the appropriate homogeneous boundary conditions, say $u_k = 0$, it follows that

$$u(X,Y) = \sum_k V_k(X) u_k(Y) \quad (19)$$

$$g(X,Y) = \sum_k G_k(X) u_k(Y) \quad (20)$$

where V_k satisfies

$$\bar{\rho} V_k + \frac{\partial^2 V_k}{\partial X^2} + \lambda_k V_k = G_k. \quad (21)$$

Then (21) is solved by Fourier transformation in X , so the solution u to (10) or (13) can be found by further matrix operations.

In the next Section, results are reported from this code using up to 64 Fourier modes and Chebyshev polynomials in each of the X and Y directions. Typically, the fractional step method with local extrapolation and tensor stress evaluation requires about 6 ms per retained mode per time step on a CDC 6600 computer. Time steps are restricted for stability to be less than the convective stability limit $\Delta x/U$.

III. LAMINAR FLOW RESULTS

Results from three laminar flow runs will be described here. The first case is flow over a very gentle wave $a \cos kx$ with $k = 4\pi \text{cm}^{-1}$, $ka = \pi/1000$ at a Reynolds number $Re_x = 4 \times 10^5$ at $x = 29 \text{ cm}$, so the boundary thickness $\delta = 6.5 \sqrt{\frac{\nu x}{U_\infty}} = 0.3 \text{ cm}$. A plot of the mapped grid used in the calculations of this flow is given in Fig. 1, where only the portion of the grid near the wall is plotted. Since the x and y scales are different, the conformal nature of the map is obscured in the plot.

In Fig. 2 we plot contours of the pressure at $t = 90$ (in nondimensional units in which $U_\infty = 1$ after the solution has converged to a steady state. The calculation was performed using 32 Fourier modes in X and 33 Chebyshev polynomials in Y . The contouring is done in the unmapped $x_1 = x$ and $x_2 = y$ coordinates. The phase of the pressure distribution lags $214^\circ \pm 5^\circ$ behind the surface displacement $a \cos x$ and the pressure coefficient $C_p = |p|_{\max} / \frac{1}{2} \rho U_\infty^2$ is

$$C_p = 1.76 \times 10^{-4}$$

at the wall. As mentioned above, the Kelvin-Helmholtz inviscid analysis predicts $C_p = 2ka = 6.3 \times 10^{-3}$. On the other hand, Benjamin's theory³ predicts a phase shift of about 217° with C_p in good agreement with our numerical simulations.

In Figs. 3 and 4 we plot contours of the velocity in the x_2 direction in the unmapped $x_1 = x$ and $x_2 = y$ coordinates. Fig. 4 is a blown-up portion of Fig. 3 that highlights the structure of the velocity v_2 near the wall. Note that the scale of the boundary layer is such that $u = 0.96U_\infty$ at $ky = 2.5$. In Fig. 5 we plot contours of u_1 in the x - y plane blown-up to highlight the wall region. In Figs. 6 and 7, we plot contours of the spanwise vorticity at $t = 90$. Again, Fig. 7 is a blown-up version of Fig. 6.

The second run is similar to the first except that the amplitude of the surface displacement is increased to

$$ka = 5.5 \times 10^{-2}$$

In Fig. 8, a contour plot is given of the pressure distribution at a nondimensional time $t = 90$. The phase of the pressure distribution now lags about 212° behind the surface while the pressure coefficient is increased to about

$$C_p = 3.3 \times 10^{-3}.$$

Benjamin's theory still predicts a phase shift of about 217° with C_p in good agreement with the value obtained by numerical simulation. In Figs. 9-11, we give contour plots in the wall region for the velocity components v_1, v_2 and spanwise vorticity, respectively, in the second run.

The third run involves the simulation of flow studied experimentally by Kachanov et al⁹. Here $k = 2.5 \text{ cm}^{-1}$ while

$$ka = (5.3 \pm 0.1) \times 10^{-2}.$$

The calculation was performed at a Reynolds number $Re_x = 1.1 \times 10^5$ to match the experiment, while $U_\infty = 530 \text{ cm/s}$ so the boundary layer thickness $\delta = 6.5 \sqrt{\frac{\nu x}{U_\infty}}$ at $x = 3.06 \text{ cm}$ is $\delta = 0.6 \text{ cm}$. Because the wavelength of the wall displacement is much longer in this run than in the second run, it turns out that the accurate simulation of this case is somewhat more difficult. Pressure perturbations due to the wall die out over a distance of order the wall wave-length which is much larger than δ so the resolution requirements in Y are more severe than for the first two cases. The results reported below were obtained using 32 Fourier modes in X and 65 Chebyshev polynomials in Y .

In Fig. 12, we plot contours of the pressure field for the third run after a steady state has been achieved. In Figs. 13 and 14, contour plots of the spanwise vorticity and velocity v_1 , respectively, near the wall are given. Kachanov et al⁹ measured the distribution of x -velocity $v_1 = u$ fluctuations at a height of about δ above the wall. They found a fluctuation amplitude Δu of

$$\frac{\Delta u_{\text{expt}}}{U_\infty} = 0.005 \pm 0.002$$

with a phase shift of $40^\circ \pm 5^\circ$ relative to the wall. The computational results give

$$\frac{\Delta u_{\text{comp}}}{U_\infty} = 0.004 \pm 0.0002$$

with a phase shift of $22^\circ \pm 5^\circ$. The magnitude of the fluctuations agree well between laboratory and computer experiment. The discrepancy in the phase shifts may be due to the sensitivity of this quantity to measurement error or due to the precise manner of measurement (in the computer experiment, the fluctuations of u were measured at constant Y while in the laboratory experiment the measurements seem to have been conducted at constant y).

IV. TURBULENT FLOW RESULTS

In this Section, we report a simulation of Kendall's experiment¹⁰ of turbulent flow over a wavy wall. An eddy viscosity model is used to simulate the effects of the turbulence. The turbulent Reynolds stresses are modelled as

$$-\overline{\rho v_1'^2} = 2\mu_T \partial v_1 / \partial x_1 - \frac{2}{3} \rho \bar{k} \quad (22)$$

$$-\overline{\rho v_1' v_2'} = \mu_T (\partial v_1 / \partial x_2 + \partial v_2 / \partial x_1) \quad (23)$$

$$-\overline{\rho v_2'^2} = 2\mu_T \partial v_2 / \partial x_2 - \frac{2}{3} \rho \bar{k} \quad (24)$$

where primes indicate turbulent fluctuations, overbars are averages, μ_T is a (spatially varying) eddy viscosity coefficient and \bar{k} is the turbulent kinetic energy.

The eddy coefficient μ_T is modelled by use of a turbulent boundary layer code¹¹ that includes effects of wall curvature and pressure gradients within a one-equation eddy-viscosity model.¹² The actual pressure measurements of the Kendall experiment were used in this boundary layer code. The turbulent kinetic energy k is modelled on the basis of flat plate data according to

$$-\overline{v_1'v_2'} / \bar{k} \approx 0.15 \quad (25)$$

A contour plot of the resulting eddy viscosity coefficient μ_T is given in Fig. 15.

Kendall's experiment is simulated using a wall displacement $a \cos(kx)$ with $a = 0.32$ cm and $k = 0.62$ cm⁻¹. The mean flow velocity in the free stream is $U = 500$ cm/s while the molecular viscosity is 0.14 cm²/s. The calculations are performed using 32 Fourier modes in X and 33 Chebyshev polynomials in Y . The flow field is initialized using the result obtained from the boundary layer code¹¹ corrected to be incompressible.

In Fig. 16, we plot contours of the velocity field v_1 in the x_1 -direction at $t = 2$ s. In Fig. 17, spanwise vorticity contours at $t = 2$ s are shown. Finally, in Figs. 18 and 19, contour plots of v_1 and v_2 , respectively, at $t=4$ s are given. It can be observed

from Fig. 18 that a small separation bubble is developing in the flow at $t=4s$. Separation seems to begin near $t=3.2s$. While no separation was observed in Kendall's experiment, it is known that the laboratory flow is close to separation. There are several possible reasons for the appearance of a separation bubble in the computation. First, the eddy viscosity distribution μ_T is fixed for all time by the turbulent boundary layer calculation. Since the mean-flow velocity has evolved significantly from the boundary layer approximation, it seems reasonable to expect that μ_T should change also. Second, the modelling of normal Reynolds stresses in (22)-(24) may be suspect. A possible remedy for both these problems would be integrating a multi-equation turbulence model at the same time as the mean flow equations. Further work is now planned in this direction. Finally, while the numerical resolution of the flow seems adequate in the wall region, there does seem to be need for more resolution at large distances from the wall. For this reason, higher resolution calculations with modified mappings in the y direction are now underway.

V. CONCLUSION

A computer code to solve the full two-dimensional Navier-Stokes equations in an arbitrary wavy geometry has been developed. The code has been run on both laminar and turbulent flows over wavy walls. The results have been compared with both available theory and experiment. For laminar flow, good agreement is achieved. For turbulent flow, the agreement is less satisfactory and may be explained by the simplified turbulence modelling used in these calculations.

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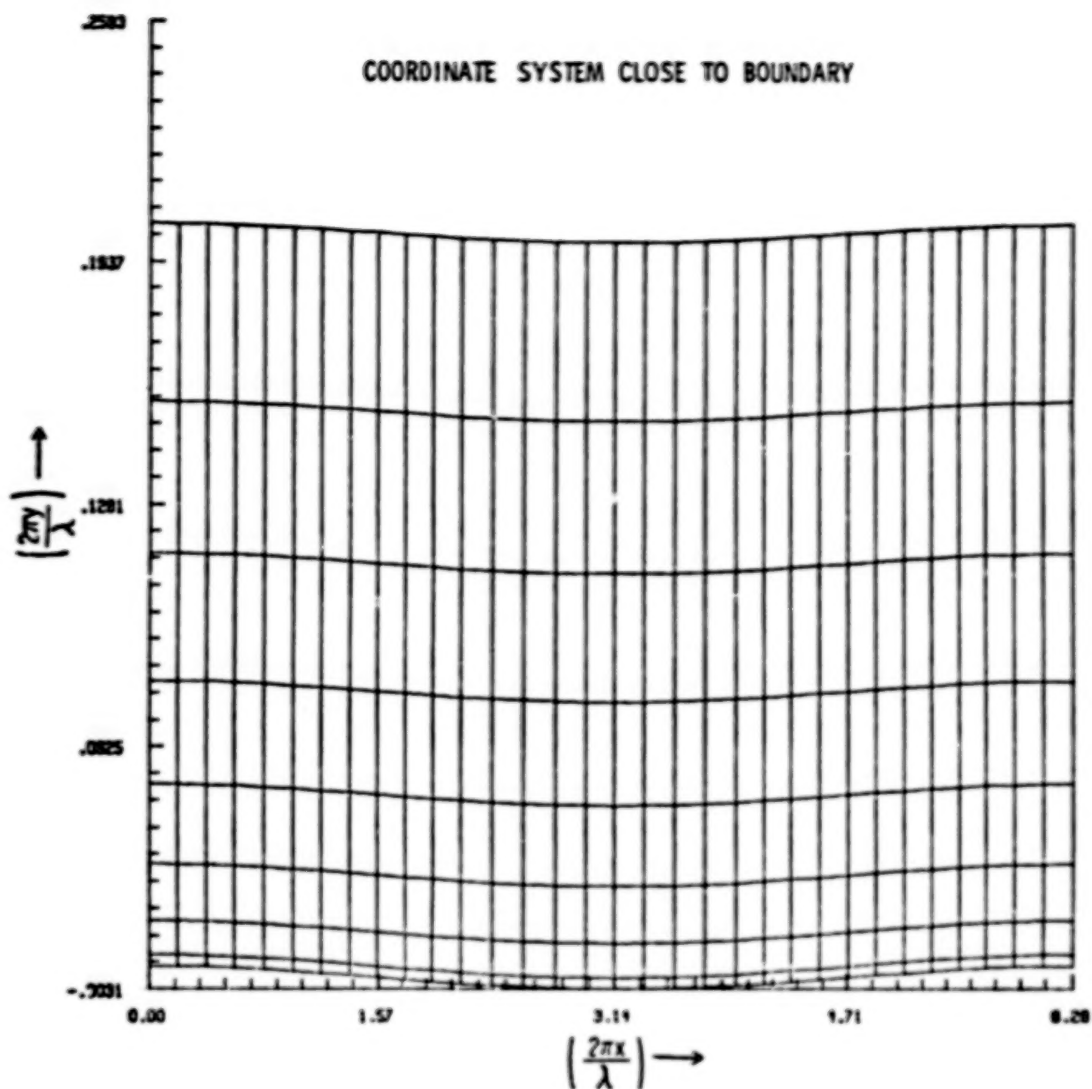


Figure 1. A plot of the coordinate system near the wall
for Run 1. Here the wall displacement is $y = a \cos kx$
with $k = 10\pi$ and $a = 10^{-4}$.

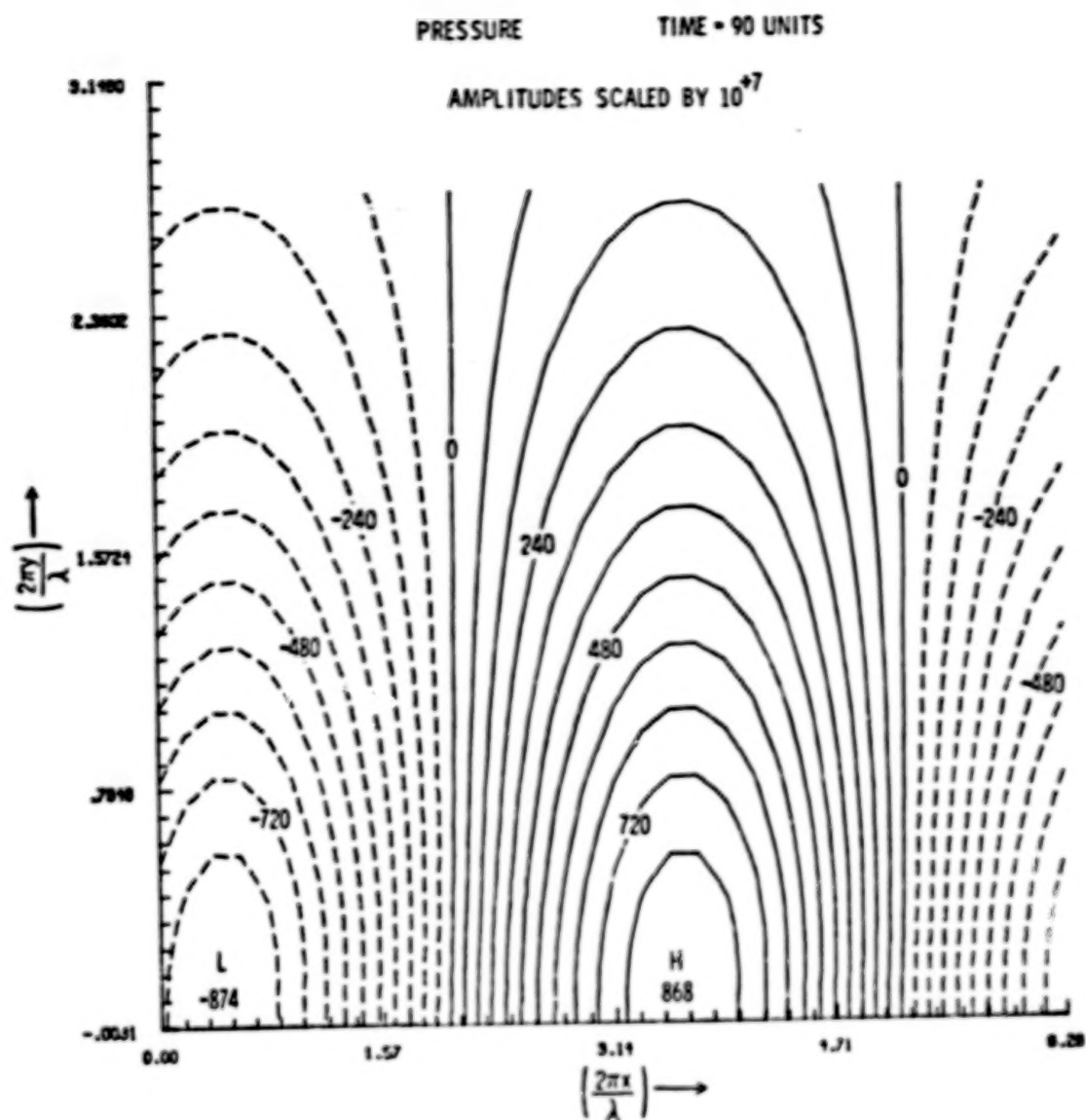


Figure 2. Contour plot of the pressure p at $t=90$ for Run 1.

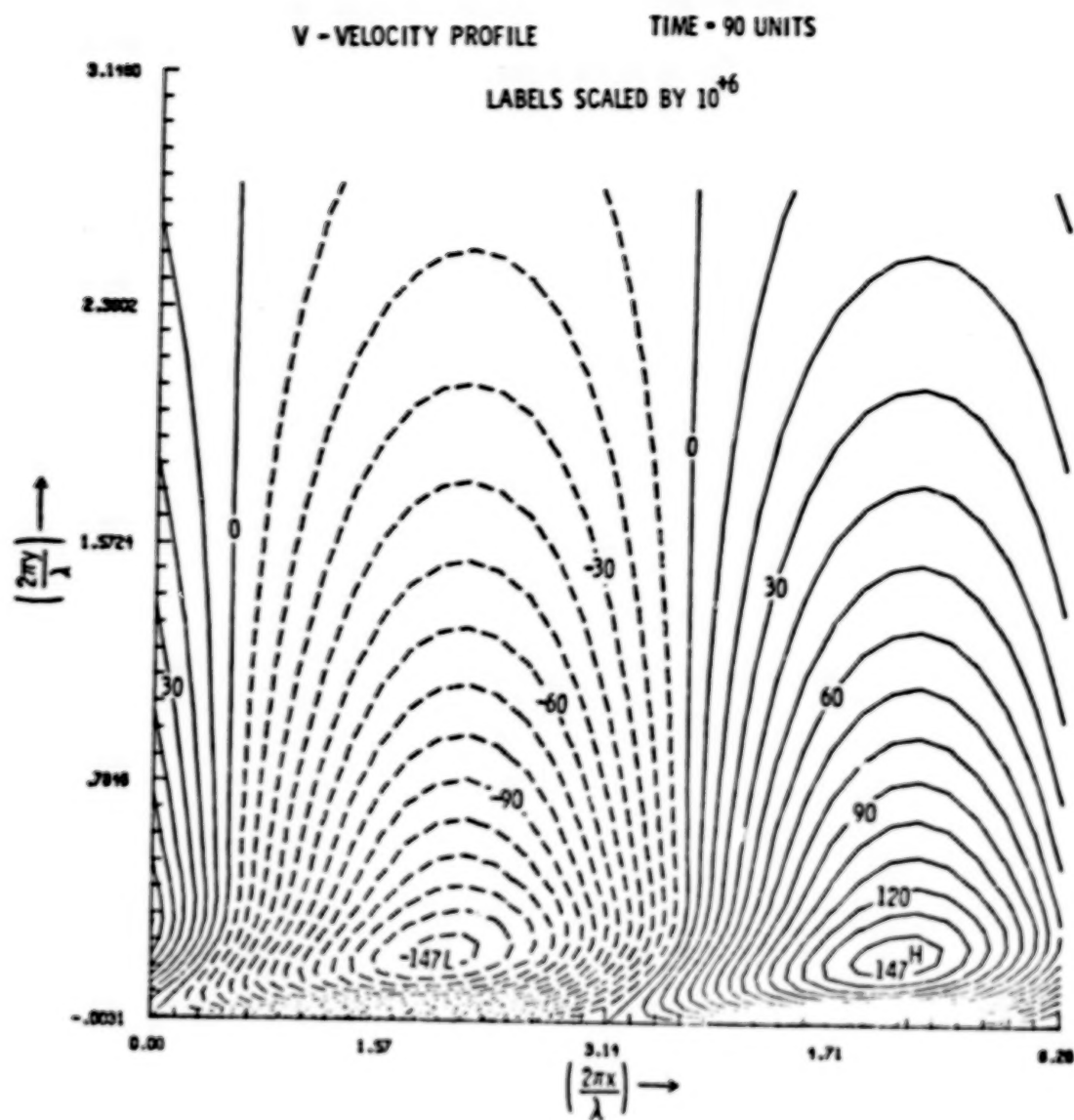


Figure 3. Contour plot of the x_2 -velocity v_2 at $t=90$ for Run 1.

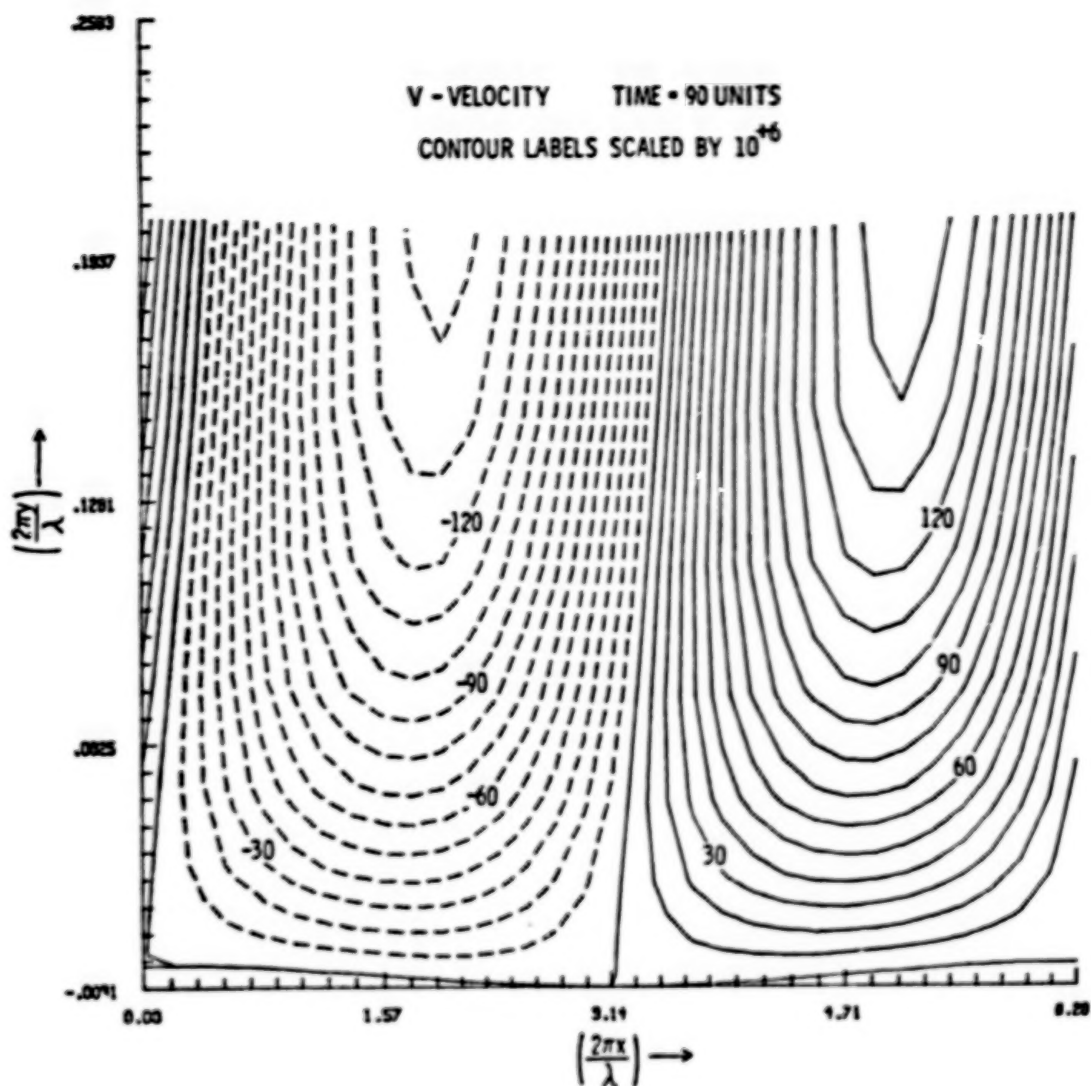


Figure 4. Same as Figure 3 except that the region near the wall is expanded.

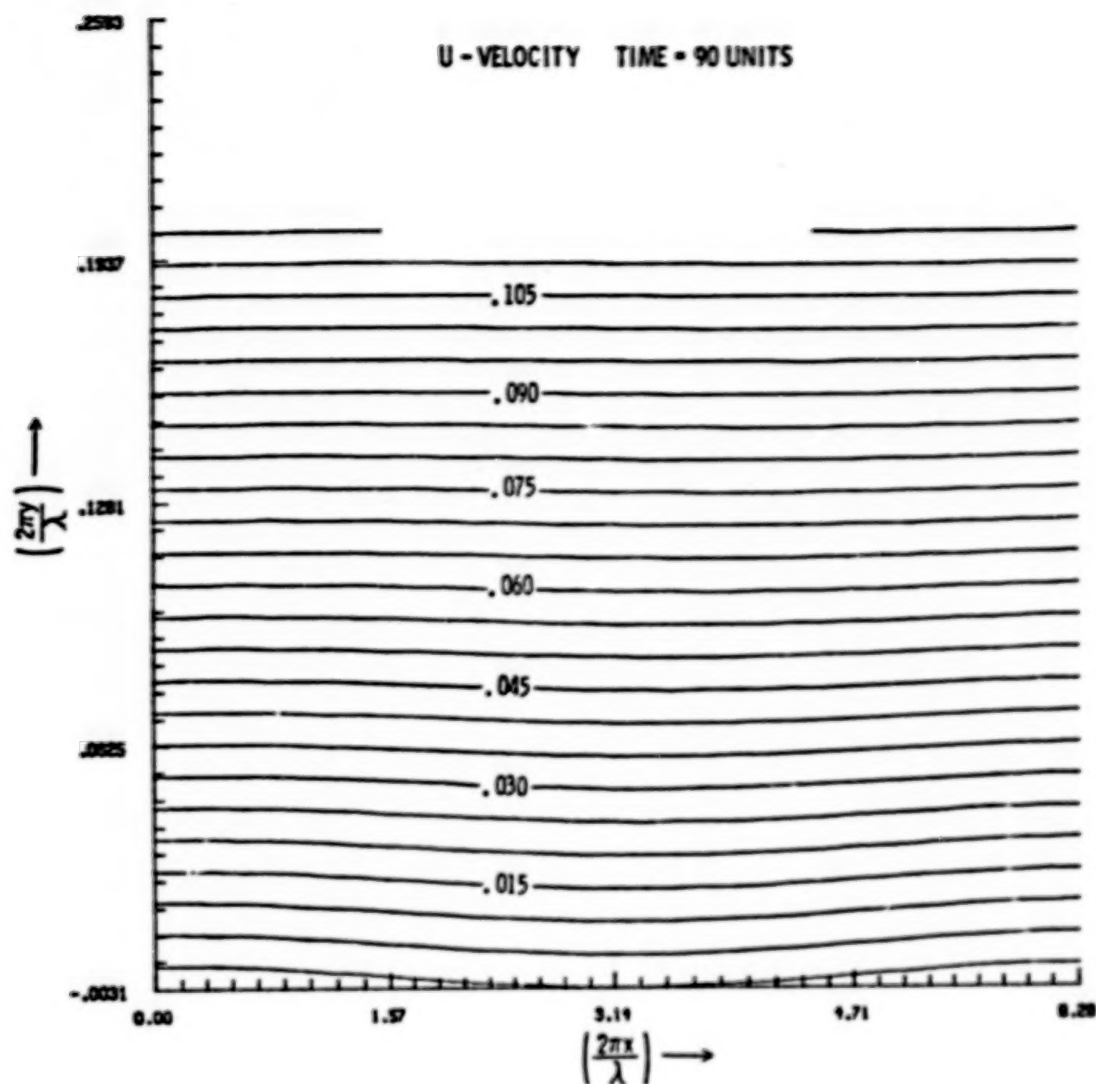


Figure 5. Contour plot of the x_1 -velocity v_1 at $t=90$ near the wall for Run 1.

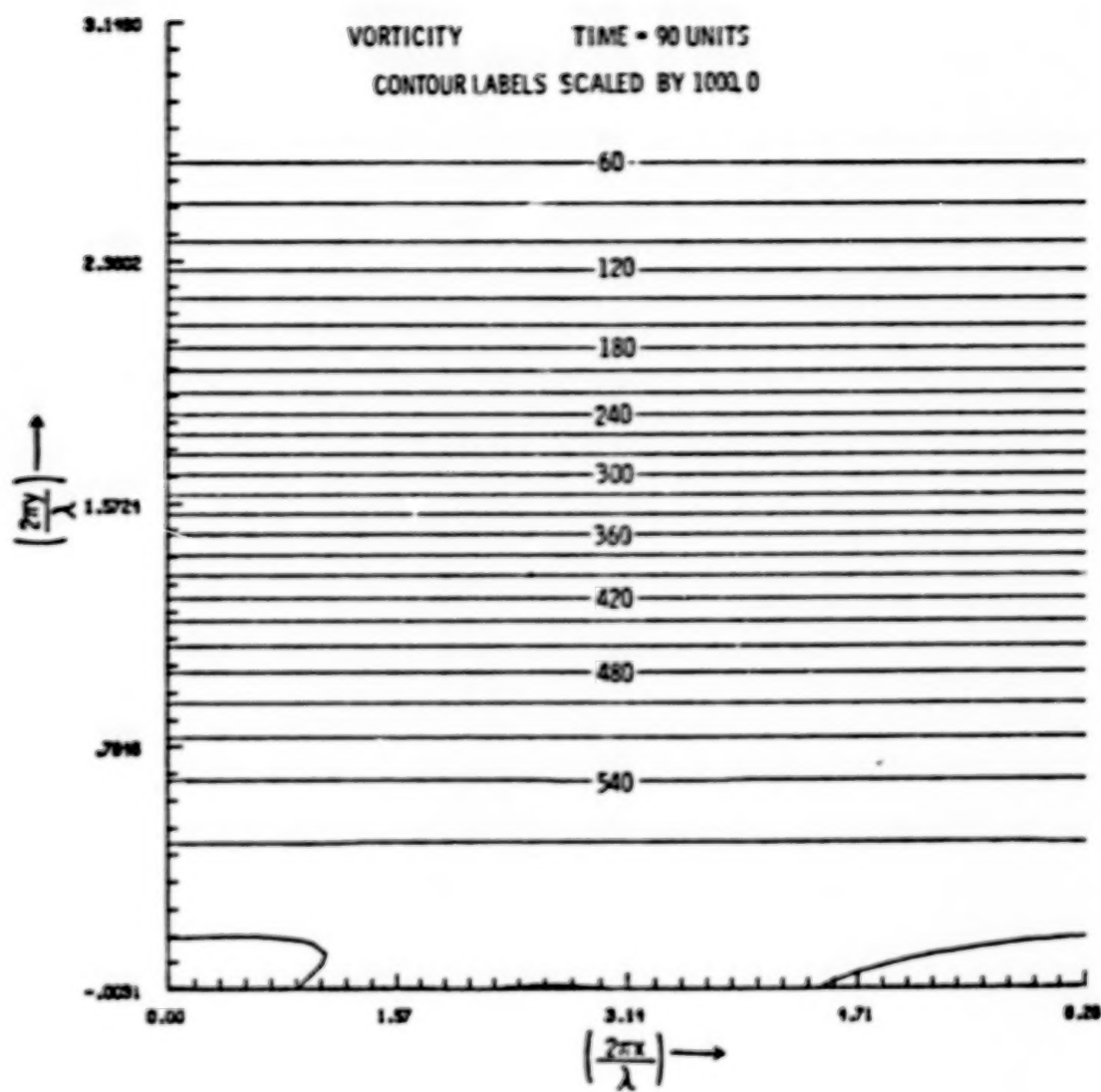


Figure 6. Contour plot of the spanwise vorticity in
Run 1 at $t=90$.

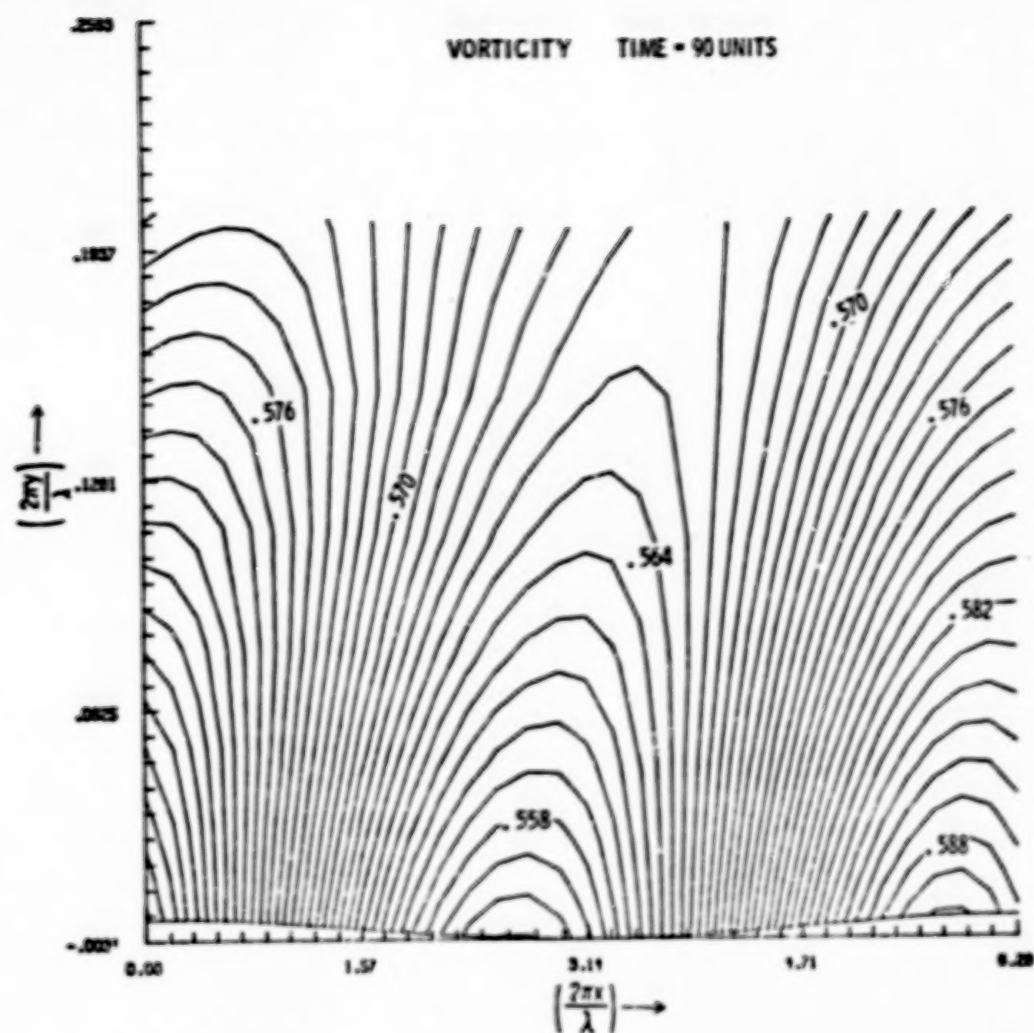


Figure 7. Same as Figure 6 except that the region near the wall is expanded.

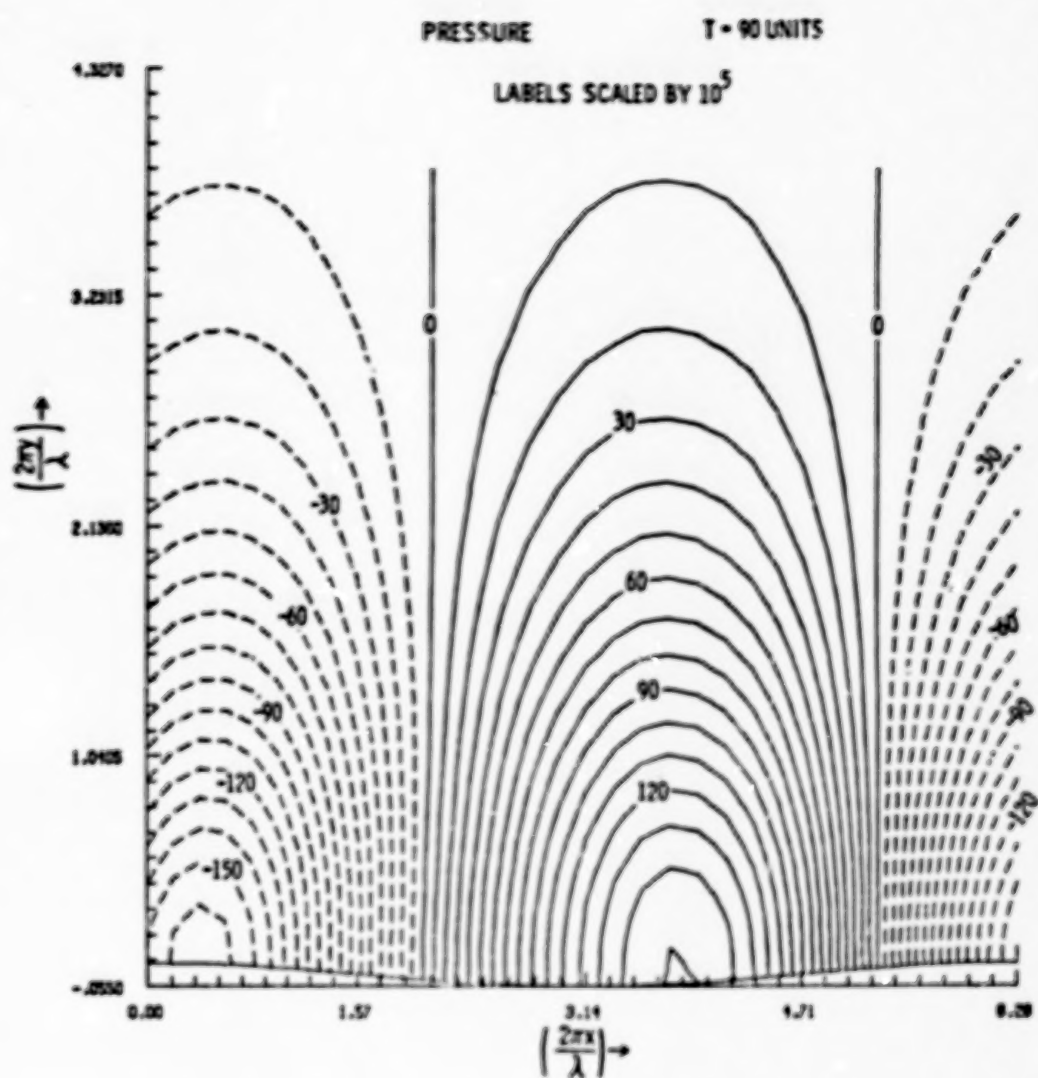


Figure 8. Contour plot of the pressure p at $t=90$ in Run 2.

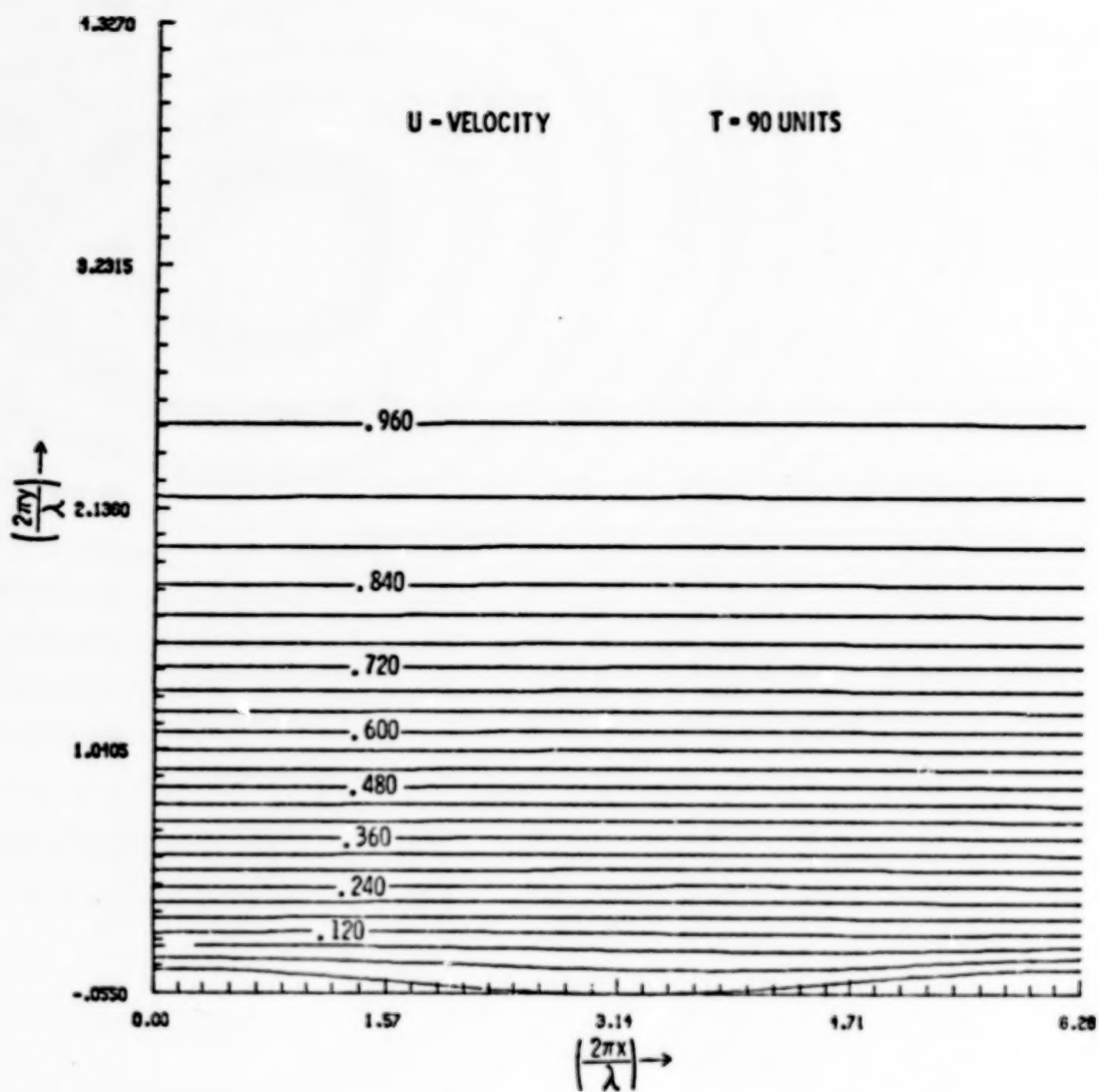


Figure 9. Contour plot of v_1 at $t=90$ in Run 2.

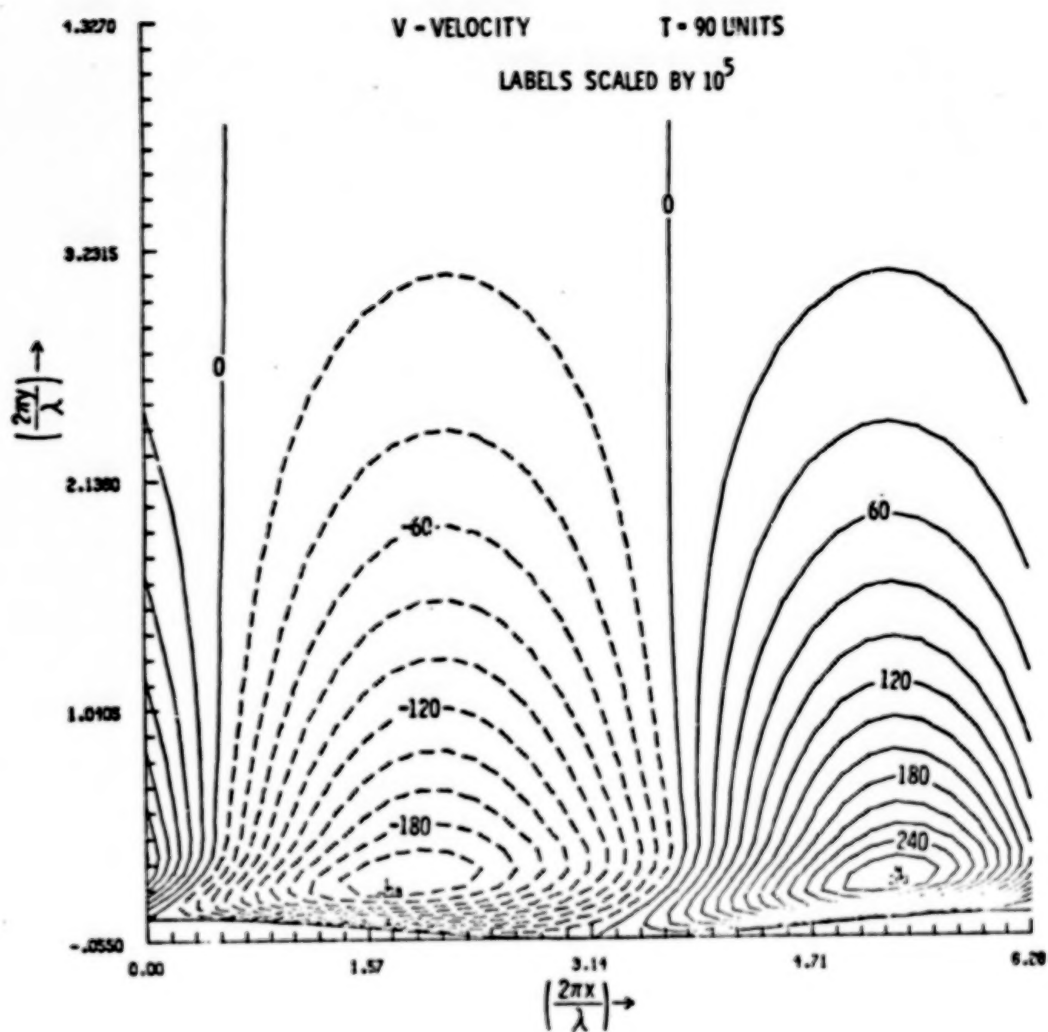


Figure 10. Contour plot of v_2 at $t=90$ in Run 2.

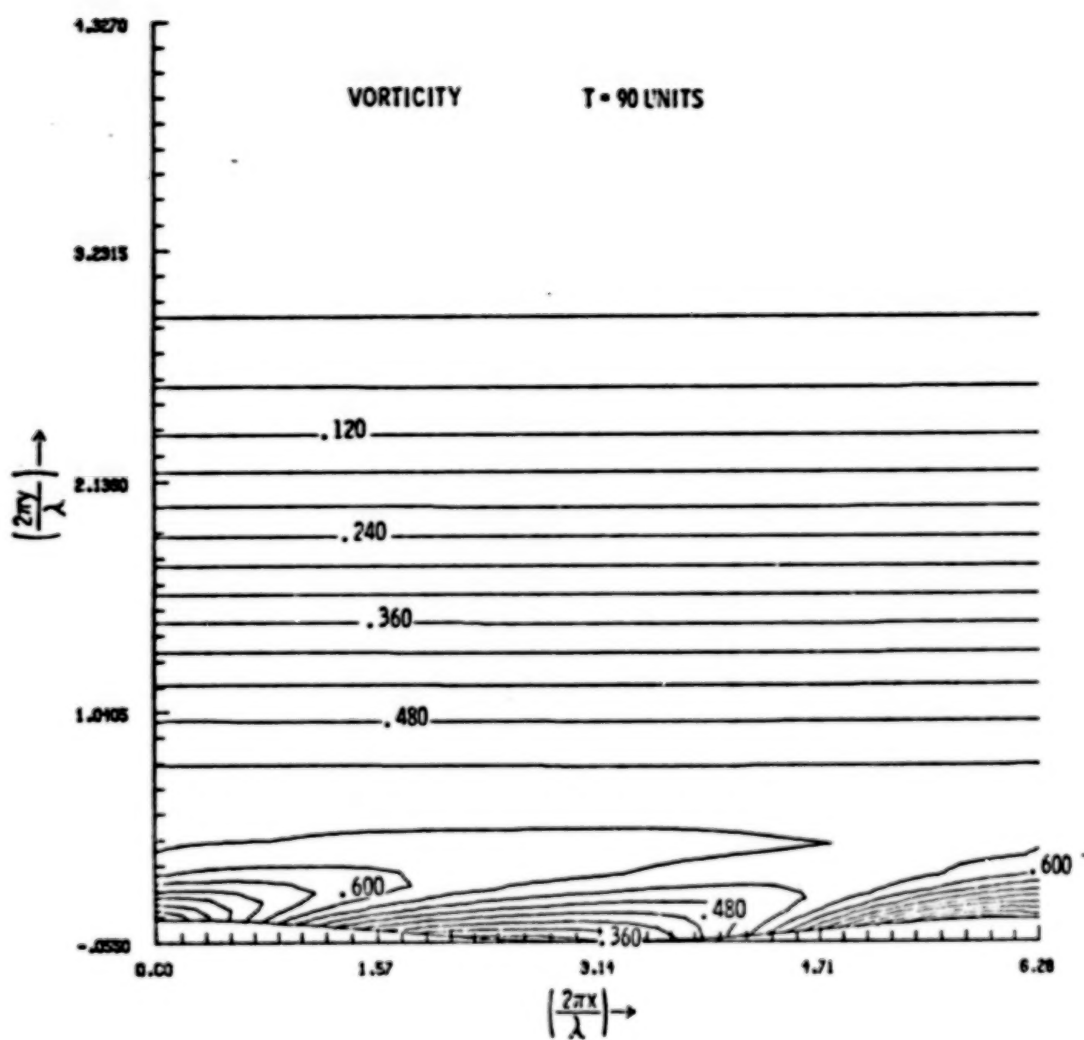


Figure 11. Contour plot of the spanwise vorticity in Run 2
at $t=90$.

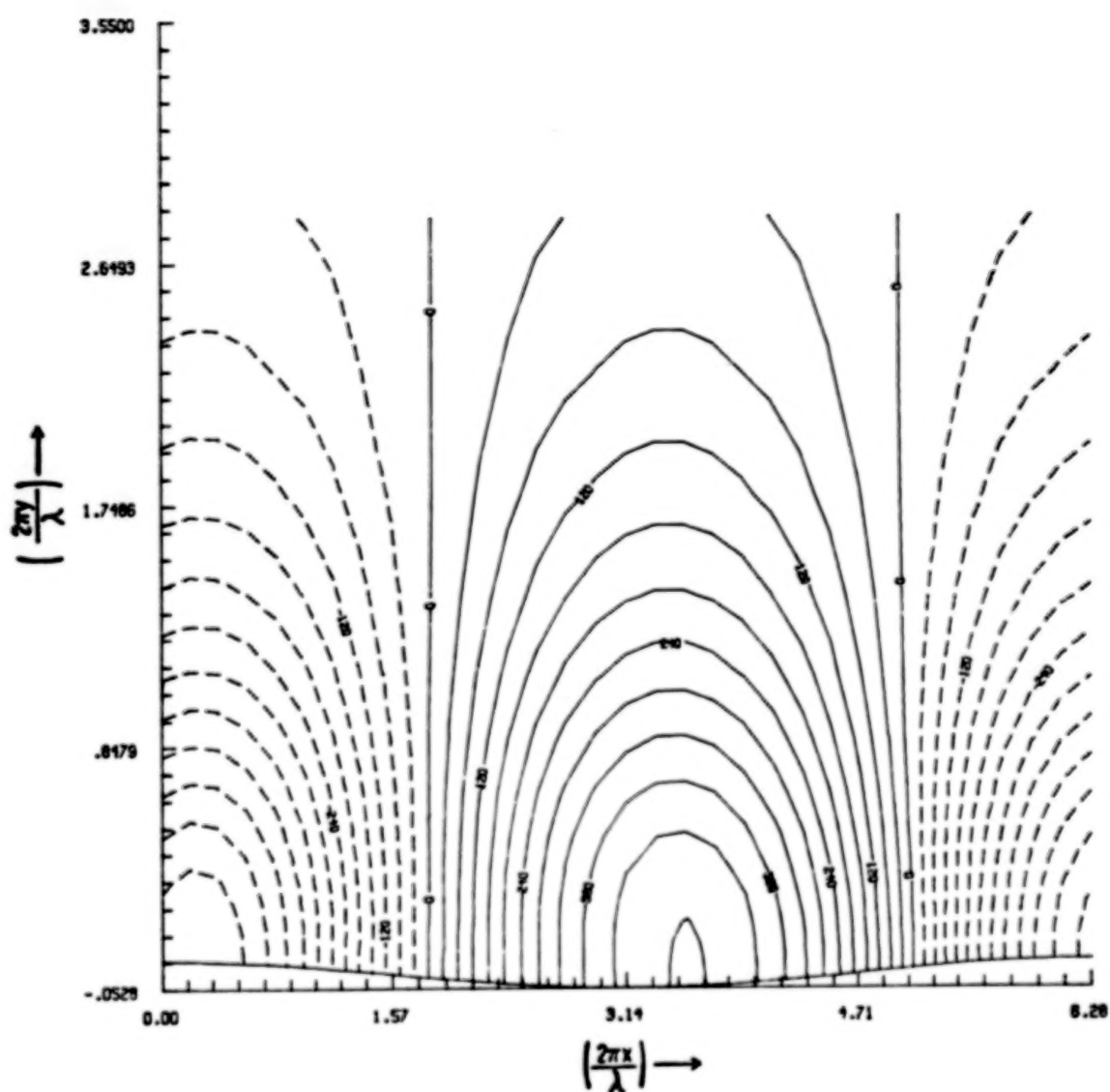


Figure 12. Contour plot of the pressure p after a steady state has been achieved in Run 3.

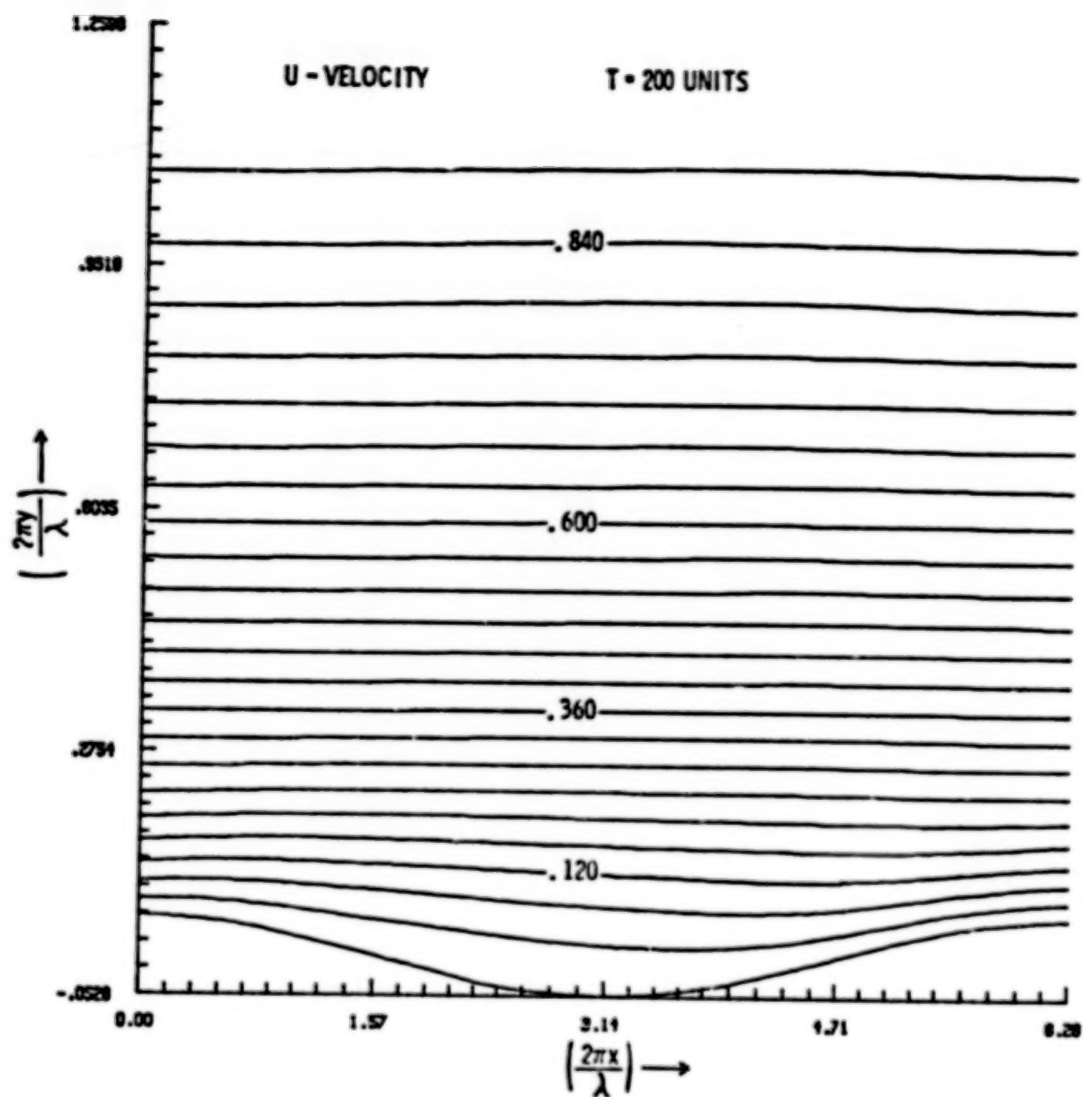


Figure 13. Contour plot of the x-velocity v_1 in Run 3 after a steady state is achieved.

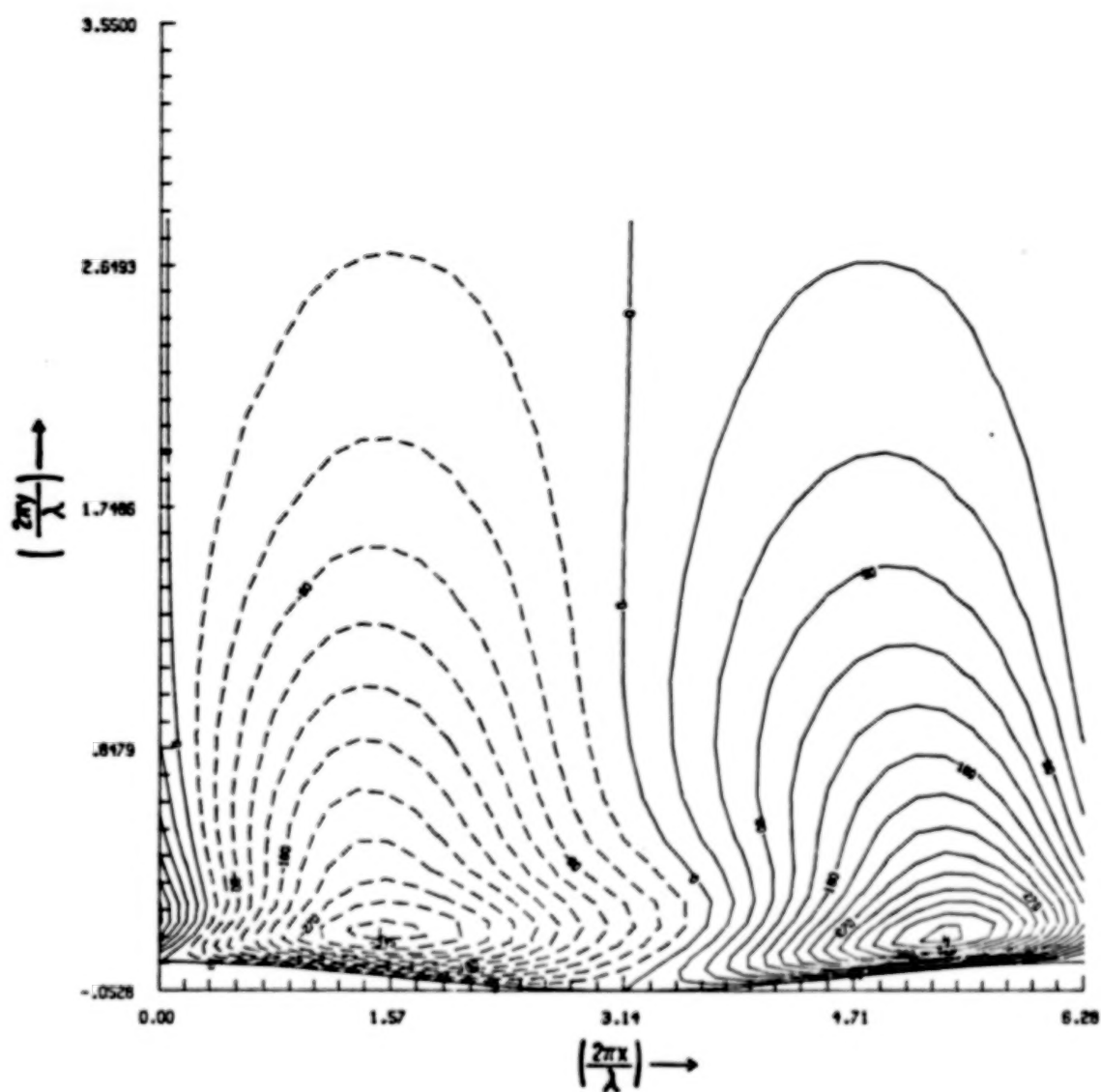


Figure 14. Contour plot of the y-velocity v_2 in Run 3 after a steady state is achieved.

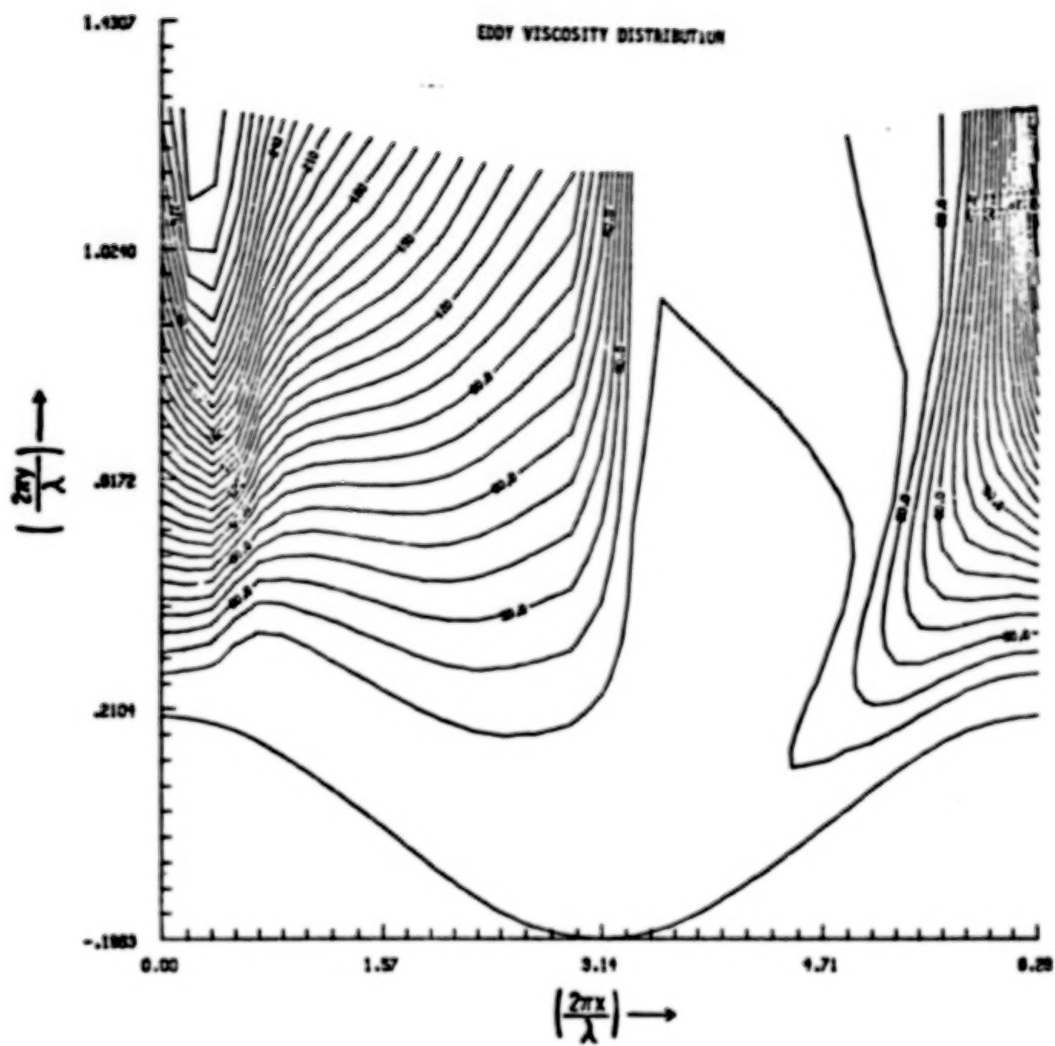


Figure 15. Contour plot of the eddy viscosity coefficient μ_T used to simulate Kendall's experiment.¹⁰

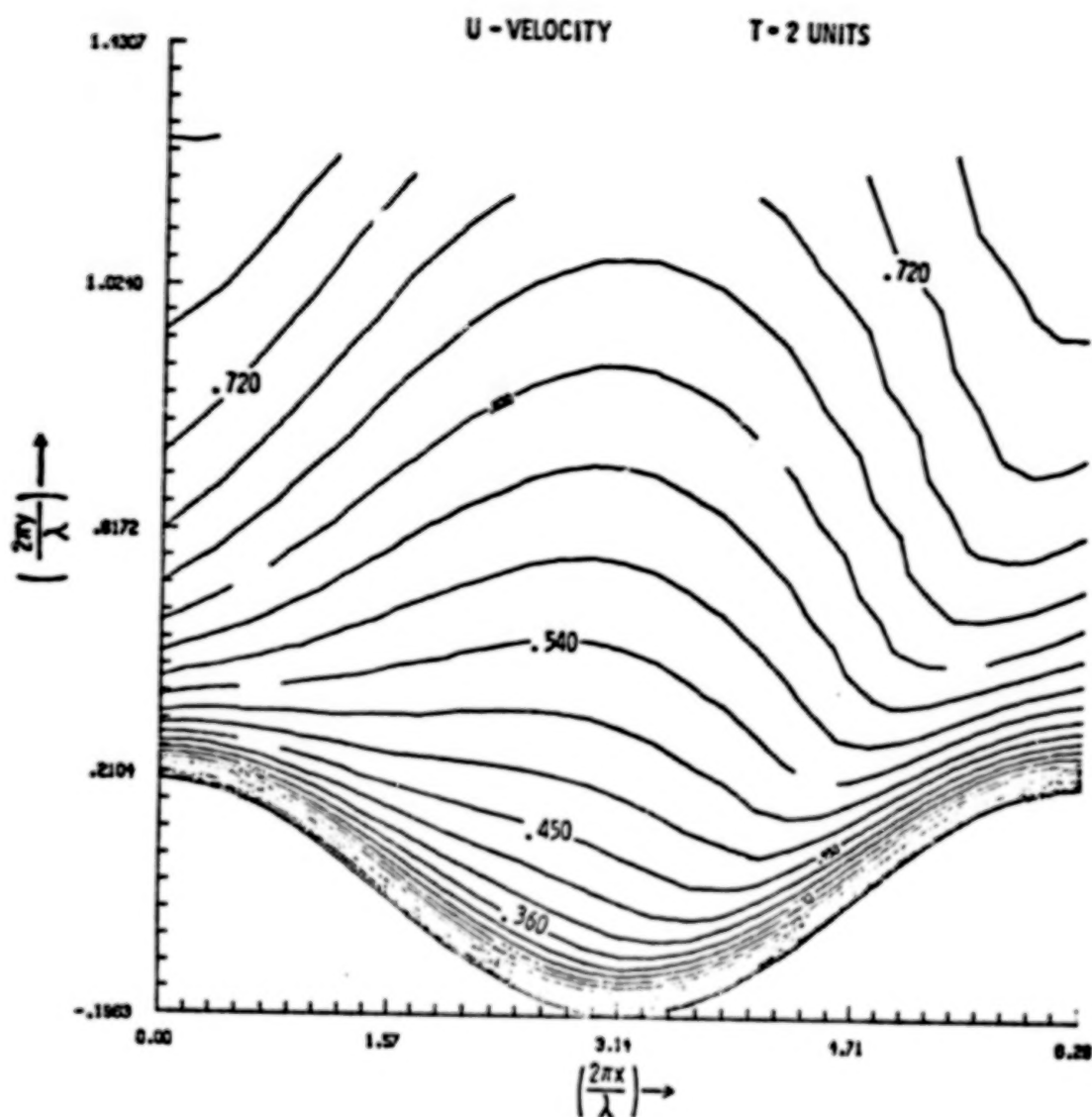


Figure 16. Contour plot of v_1 at $t=2s$ in the simulation of Kendall's experiment.

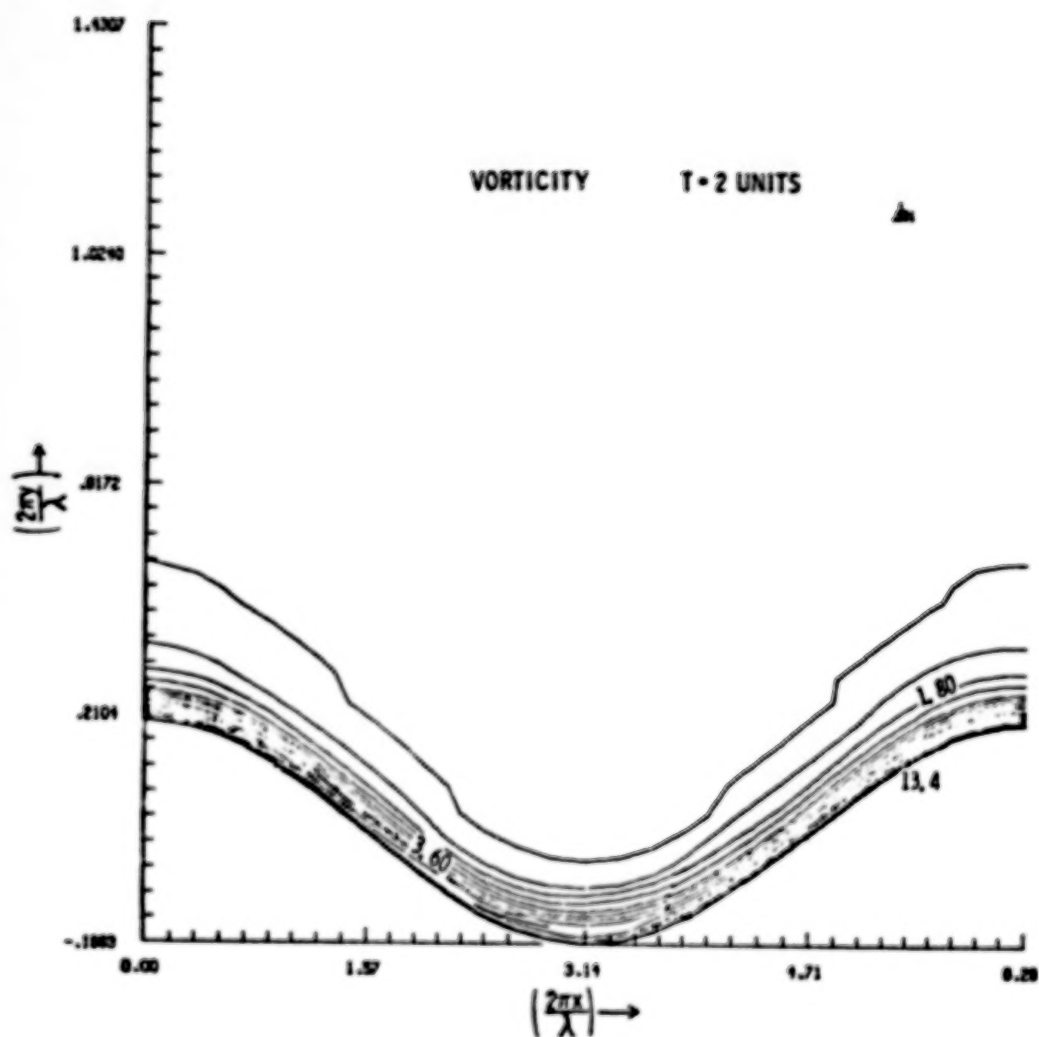


Figure 17. Same as Figure 16 except for the spanwise vorticity.

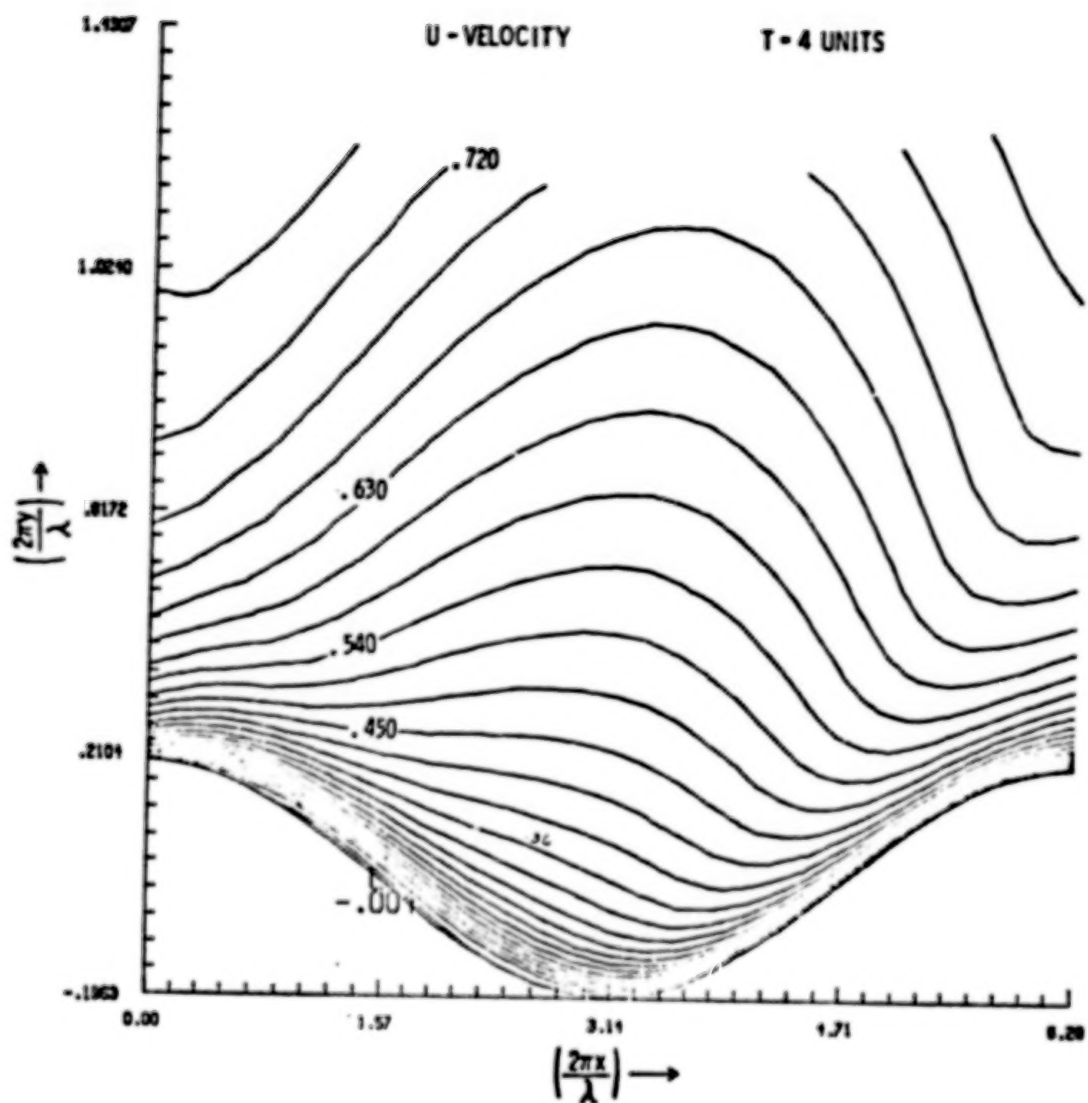


Figure 18. Contour plot of v_1 at $t=4s$ in the simulation of Kendall's experiment.

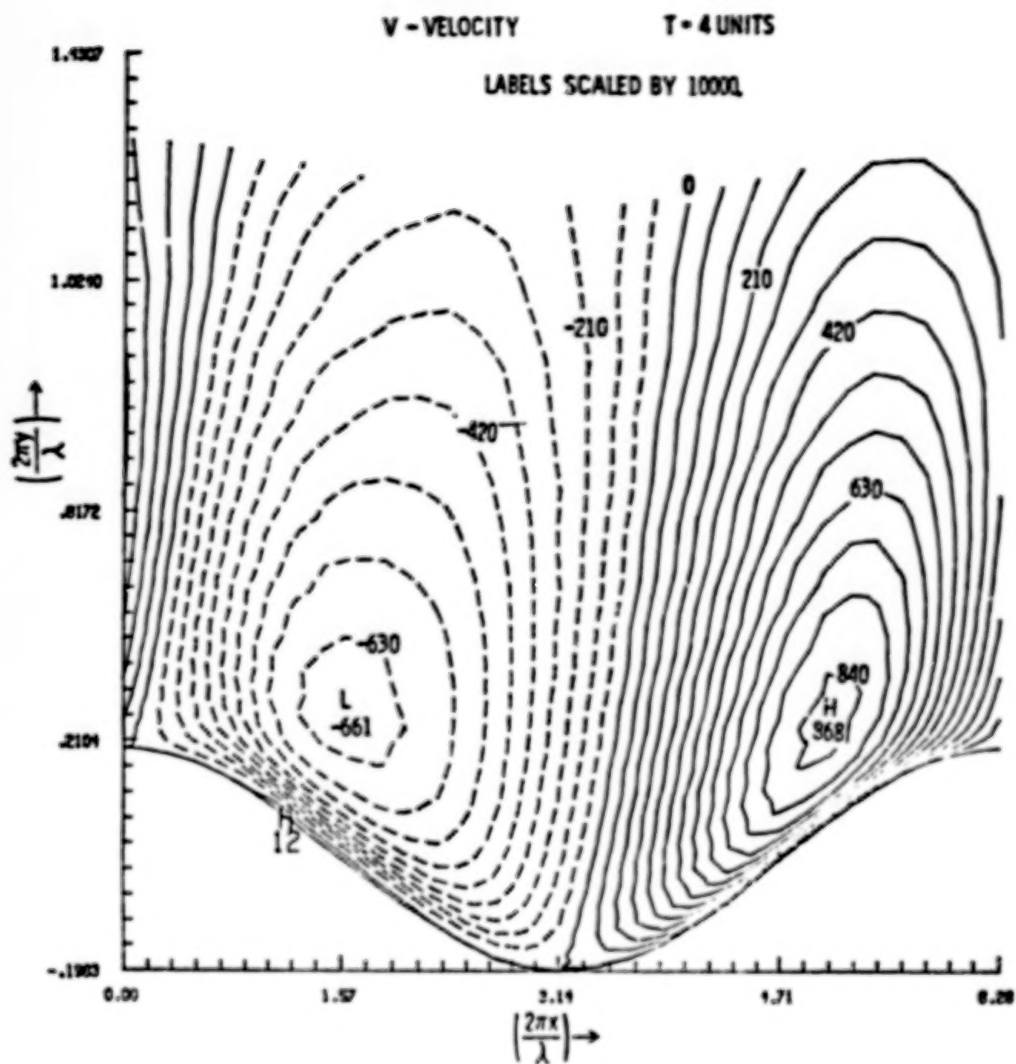
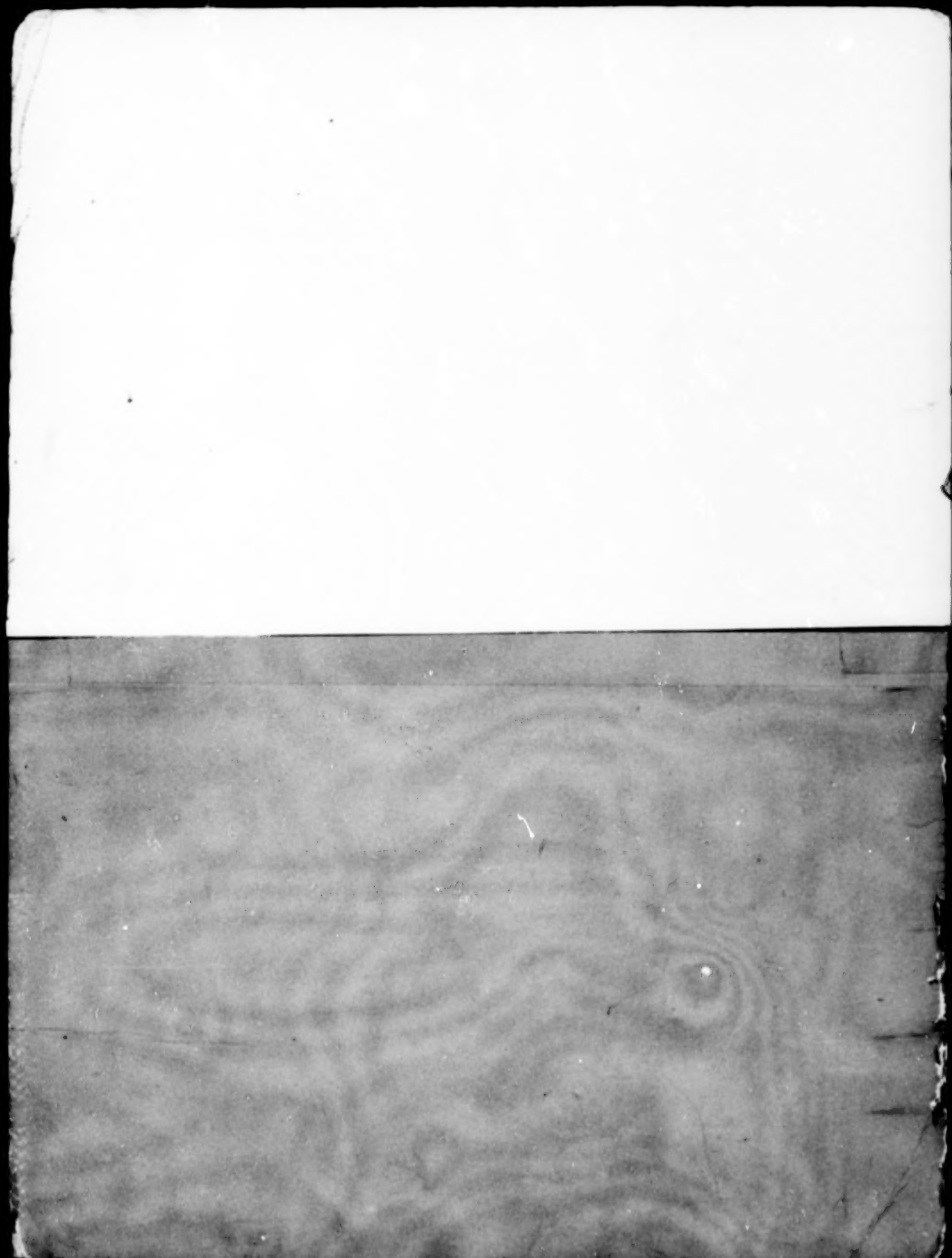


Figure 19. Same as Figure 18 except for v_2 .

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